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LETTER TO THE EDITOR

On localized solitonic solutions of a (2+1)-dimensional sine-Gordon system

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Abstract. Exponentially decaying solutions of an integrable sine-Gordon system in 2+1 dimensions are presented. A simple superposition formula is given.

It has recently been shown that the (2+1)-dimensional integrable extension of the classical sine-Gordon equation found by Konopelchenko and Rogers [1]

$$\left(\frac{\theta_{tx}}{\sin\theta}\right)_{x} - \left(\frac{\theta_{ty}}{\sin\theta}\right)_{y} + \frac{\theta_{x}\tilde{\theta}_{y} - \theta_{y}\tilde{\theta}_{x}}{\sin^{2}\theta} = 0$$
$$\left(\frac{\tilde{\theta}_{x}}{\sin\theta}\right)_{x} - \left(\frac{\tilde{\theta}_{y}}{\sin\theta}\right)_{y} + \frac{\theta_{x}\theta_{ty} - \theta_{y}\theta_{tx}}{\sin^{2}\theta} = 0$$

is amenable to an extended version of a Darboux-Levi transformation [2]. The simplest solutions that have been obtained by this approach have the following form:

$$\theta = 4 \tan^{-1} \left(\frac{|D^+|}{|D^-|} \right)$$

where

$$D^{\pm} := \begin{pmatrix} 0 & e^{-\alpha_{1}} & \dots & e^{-\alpha_{N}} \\ e^{\pm \alpha_{1}} & & \\ \vdots & M \\ e^{\pm \alpha_{N}} & & \end{pmatrix}$$
$$M_{ij} := 2 \frac{\lambda_{j}}{\lambda_{i} + \lambda_{j}} e^{\alpha_{i} + \alpha_{j}} + 2 \frac{\mu_{j}}{\mu_{i} + \mu_{j}} \cdot e^{-\alpha_{i} - \alpha_{j}}$$

and

$$\alpha_i \coloneqq \frac{1}{2} (\lambda_i - \mu_i) x - \frac{1}{2} (\lambda_i + \mu_i) y + (\lambda_0 \lambda_i^{-1} - \mu_0 \mu_i^{-1}) t + \nu_i$$

with arbitrary real constants λ_i , μ_i , λ_0 , μ_0 and complex constants ν_i the imaginary part of which being half-integer multiplies of π .

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In order to construct localized solutions we now restrict the parameters to $\mu_i = -\lambda_i$ for $i = 1, ..., N_0$ and $\mu_i = \lambda_i$ for $i = N_0 + 1, ..., N$, that is

$$\alpha_i = \lambda_i x + (\lambda_0 + \mu_0) \lambda_i^{-1} t + \nu_i \qquad i = 1, \dots, N_0$$

$$\alpha_i = -\lambda_i y + (\lambda_0 - \mu_0) \lambda_i^{-1} t + \nu_i \qquad i = N_0 + 1, \dots, N_0$$

For N = 2, $N_0 = 1$ we get the solution

$$\theta = 4 \tan^{-1} \left(\frac{\cosh(\alpha_1 + \alpha_2)}{\sinh(\alpha_1 - \alpha_2)} \right)$$

where multiplicative factors have been absorbed into ν_1 and ν_2 .

The associated quantity

$$\theta_x \theta_y = 16 \frac{\lambda_1 \lambda_2}{\cosh 2\alpha_1 \cosh 2\alpha_2} \tag{1}$$

provides a localized object since it decays to zero exponentially as $(x, y) \rightarrow \infty$ in any direction and has the amplitude $\pm 16\lambda_1\lambda_2$ depending on the imaginary parts of ν_1 and ν_2 . Interestingly, it can be regarded as a nonlinear superposition of two one-dimensional (anti)kinks for positive amplitude and a kink and an antikink for negative amplitude moving in the x-direction and y-direction respectively, that is

$$\theta_x \theta_y = a_x b_y$$

where a and b are the well known (anti)kink solutions of two copies of the sine-Gordon equation, namely:

$$a_{xt} = 4(\mu_0 + \lambda_0) \sin a$$
$$b_{yt} = 4(\mu_0 - \lambda_0) \sin b.$$

In figure 1 the kink-kink solution (1) is plotted at a fixed time t.

For N = 3, $N_0 = 2$ it may be shown that the solution decomposes asymptotically into two solutions as described by (1) provided the quantities $(c_2 + c_{32} + c_{13}) \times (c_{21} + c_{32} + c_{13})$ and $(c_{21} + c_{32}^{-1} + c_{13}^{-1})c_2c_{32}c_{13}$ are negative, where $c_{ij} := (\lambda_i - \lambda_j)/(\lambda_i + \lambda_j)$. Asymptotic analysis shows that they behave like one-dimensional solitons, that is, the only indication that an interaction has occurred is that the two objects are phaseshifted. The amplitudes remain unchanged.



Figure 1. The kink-kink solution (1) plotted at fixed time t.



t=-1/2

t=3/2



t=0



t=2









Figure 2. Interaction of a kink-kink with a kink-kink at various times.

In figures 2 and 3 snapshots are shown at various times during the interaction of a kink-kink with a kink-kink and kink-antikink respectively.

A more detailed study of these solutions is currently being undertaken.



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